## CS 3630

Inverse Kinematics: Planar Arms


## Assigning Link Frames



- Frame $n$ is the end-effector frame. It can be attached to link $n$ in any manner that is convenient.
- In this case, $n=2$, so Frame 2 is the end-effector frame.


## The Forward Kinematic Map



Once we have coordinate frames for each link:

- Determine $T_{i}^{i-1}$ for adjacent links as a function of $q_{i}$
- The forward kinematic map is given by: $T_{n}^{0}\left(q_{1} \ldots q_{n}\right)=T_{1}^{0}\left(q_{1}\right) \ldots T_{n}^{n-1}\left(q_{n}\right)$


## The Forward Kinematic Map

- The forward kinematic map gives the position and orientation of the end-effector frame as a function of the joint variables:

$$
T_{n}^{0}=F\left(q_{1}, \ldots, q_{n}\right)
$$

- For the two-link planar arm, we have


$$
\begin{aligned}
T_{2}^{0} & =\left[\begin{array}{ccc}
\cos \theta_{1} & -\sin \theta_{1} & a_{1} \cos \theta_{1} \\
\sin \theta_{1} & \cos \theta_{1} & a_{1} \sin \theta_{1} \\
0 & 0 & 1
\end{array}\right]\left[\begin{array}{ccc}
\cos \theta_{2} & -\sin \theta_{2} & a_{2} \cos \theta_{2} \\
\sin \theta_{2} & \cos \theta_{2} & a_{2} \sin \theta_{2} \\
0 & 0 & 1
\end{array}\right] \\
& =\left[\begin{array}{ccc}
\cos \left(\theta_{1}+\theta_{2}\right) & -\sin \left(\theta_{1}+\theta_{2}\right) & a_{1} \cos \theta_{1}+a_{2} \cos \left(\theta_{1}+\theta_{2}\right) \\
\sin \left(\theta_{1}+\theta_{2}\right) & \cos \left(\theta_{1}+\theta_{2}\right) & a_{1} \sin \theta_{1}+a_{2} \sin \left(\theta_{1}+\theta_{2}\right) \\
0 & 0 & 1
\end{array}\right]
\end{aligned}
$$

## Inverse Kinematics

## The General Inverse Kinematics Problem:

Given the forward kinematic map: $T_{n}^{0}=F\left(q_{1}, \ldots, q_{n}\right)$
Solve for $q_{1}, \ldots, q_{n}$ to achieve a desired $T^{d}$
i.e., find $q_{1}^{d}, \ldots, q_{n}^{d}$ such that $F\left(q_{1}^{d}, \ldots, q_{n}^{d}\right)=T^{d}$

Why is this difficult?

- In general, $F\left(q_{1}, \ldots, q_{n}\right)$ will be nonlinear. Solving nonlinear equations is hard.
- Further, for a general $F\left(q_{1}, \ldots, q_{n}\right)$ we don't know
- Does a solution to $F\left(q_{1}, \ldots, q_{n}\right)=\mathrm{T}^{\mathrm{d}}$ exist?
- If a solution exists, is it unique?


## The Inverse Kinematic Solution

For the two-link arm, typically the goal is to place the end-effector at a desired location.

- Denote the coordinates of the origin of frame 2 by $o_{2, x}, o_{2, y}$.
- Solve for $\theta_{1}$ and $\theta_{2}$ such that

$$
\begin{gathered}
o_{2, x}=a_{1} \cos \theta_{1}+a_{2} \cos \left(\theta_{1}+\theta_{2}\right) \\
o_{2, y}=a_{1} \sin \theta_{1}+a_{2} \sin \left(\theta_{1}+\theta_{2}\right)
\end{gathered}
$$

- Recall that $a_{1}$ and $a_{2}$ are constants defined by the mechanical structure of the arm.
- This is a nonlinear set of equations in $\theta_{1}$ and $\theta_{2}$--- and nonlinear equations can be very difficult to solve!

$$
T_{2}^{0}=\left[\begin{array}{ccc}
\cos \left(\theta_{1}+\theta_{2}\right) & -\sin \left(\theta_{1}+\theta_{2}\right) & a_{1} \cos \theta_{1}+a_{2} \cos \left(\theta_{1}+\theta_{2}\right) \\
\sin \left(\theta_{1}+\theta_{2}\right) & \cos \left(\theta_{1}+\theta_{2}\right) & a_{1} \sin \theta_{1}+a_{2} \sin \left(\theta_{1}+\theta_{2}\right) \\
0 & 0 & 1
\end{array}\right]=\left[\begin{array}{ccc}
* & * & o_{2, x} \\
* & * & o_{2, y} \\
0 & 0 & 1
\end{array}\right]
$$

NOTE:

- We don't care about the orientation for this problem.
- In fact, we can't choose the orientation if we also choose $o_{2, x}, o_{2, y}$.


## Geometric Methods (closed-form solns)

- For some manipulators, it is possible to use fairly simple trigonometry to solve the inverse kinematics problem.
- Any two adjacent links are coplanar (any two intersecting lines are coplanar).
- The origins of frames $i-1, i$, and $i+1$ define a triangle.
- Simple, trigonometry in the plane might just get the job done!



## Solving for $\theta_{2}$



$>$ With this set of equations, we can solve for $\theta_{2}$ using simple solutions to closed-form equations.
> We never need to solve a nonlinear system!

Denote the coordinates of the origin of frame 2 by $o_{2, x}, o_{2, y}$.

The Law of Cosines:

$$
r^{2}=a_{1}^{2}+a_{2}^{2}-2 a_{1} a_{2} \cos \alpha
$$

Define:

$$
D \stackrel{\text { def }}{=} \frac{a_{1}^{2}+a_{2}^{2}-r^{2}}{2 a_{1} a_{2}}=\cos \alpha
$$

Then $\sin \alpha= \pm \sqrt{1-D^{2}}$

Finally,

$$
\alpha=\tan ^{-1} \frac{ \pm \sqrt{1-D^{2}}}{D}
$$

## What about existence and uniqueness?

Does a solution always exist for $\alpha=\tan ^{-1} \frac{ \pm \sqrt{1-D^{2}}}{D}$ ?
No solution exists if $D^{2}>1$ :

$$
\begin{gathered}
D^{2}=\left(\frac{a_{1}^{2}+a_{2}^{2}-r^{2}}{2 a_{1} a_{2}}\right)^{2} \leq 1 \\
a_{1}^{2}+a_{2}^{2}-r^{2} \leq \pm 2 a_{1} a_{2} \\
a_{1}^{2} \pm 2 a_{1} a_{2}+a_{2}^{2} \leq r^{2} \\
\left(a_{1} \pm a_{2}\right)^{2} \leq r^{2} \\
\left|a_{1} \pm a_{2}\right| \leq r
\end{gathered}
$$

In this case, a solution exists!

## What about existence and uniqueness?

 Is the solution unique for $\alpha=\tan ^{-1} \frac{ \pm \sqrt{1-D^{2}}}{D}$ ?Clearly, the solution is not unique, since we may choose either square root!
The second solution uses $\alpha=\tan ^{-1} \frac{-\sqrt{1-D^{2}}}{D}$ which results in an "elbow UP" configuration.


## Degenerate Solutions



Arm "folds back" on itself

$$
a_{1}-a_{2}=r
$$

## Solving for $\theta_{1}$



Denote the coordinates of the origin of frame 2 by $o_{2, x}, o_{2, y}$.

$$
\begin{aligned}
& \theta_{1}=\beta-\gamma \\
& \beta=\tan ^{-1} \frac{o_{2, y}}{o_{2, x}}
\end{aligned}
$$

The Law of Cosines, this time for $\gamma$ :

$$
a_{2}^{2}=a_{1}^{2}+r^{2}-2 a_{1} r \cos \gamma
$$

$>$ With this set of equations, we can solve for $\theta_{1}$ using simple solutions to closed-form equations.

- We never need to solve a nonlinear system!


## Elbow up is left as an exercise for you!

## Position and Orientation

Suppose we wish to position the end effector frame at a specific position, and with a specific orientation.

- We can parameterize the end effector frame by $\left(X_{e}, Y_{e}, \phi\right)$
- $\left(X_{e}, Y_{e}\right)$ give the position of the origin of the frame
- $\phi$ gives the orientation of the frame:

$$
T_{2}^{0}=\left[\begin{array}{ccc}
\cos \phi & -\sin \phi & X_{e} \\
\sin \phi & \cos \phi & Y_{e} \\
0 & 0 & 1
\end{array}\right]
$$

- We can't do this with a two-link arm.
- Intuitively, we have three inputs ( $\theta_{1}$ and $\theta_{2}$ ) and three outputs.
- Our solution to the two-link arm shows that once we choose $\theta_{1}$ and $\theta_{2}$ the orientation of the end-effector frame is fully determined.


## $>$ Add another link!

## Three-Link Planar Arm

$$
T_{2}^{0}=\left[\begin{array}{ccc}
C_{123} & -S_{123} & a_{1} C_{1}+a_{2} C_{12}+a_{3} C_{123} \\
S_{123} & C_{123} & a_{1} S_{1}+a_{2} S_{12}+a_{3} S_{123} \\
0 & 0 & 1
\end{array}\right]
$$

$$
C_{123}=\cos \left(\theta_{1}+\theta_{2}+\theta_{3}\right), \text { etc. }
$$

- Given $\left(\boldsymbol{X}_{e}, \boldsymbol{Y}_{e}, \boldsymbol{\phi}\right)$, we know everything about the position and orientation of link 3!
- After all, link 3 is a rigid body, with a rigidly attached Frame 3.
- If we know the position and orientation of Frame 3, we know everything about Link 3!


## Three-Link Planar Arm

Using simple trigonometry we obtain

$$
\begin{aligned}
o_{2, x}+a_{3} \cos \phi & =X_{e} \\
o_{2, y}+a_{3} \sin \phi & =Y_{e}
\end{aligned}
$$

And from this we immediately conclude

$$
\begin{gathered}
o_{2, x}=X_{e}-a_{3} \cos \phi \\
o_{2, y}=y_{e}-a_{3} \sin \phi
\end{gathered}
$$

And now, it's a 2-link arm problem


Link 1
As before, there are 0, 1, or two solutions (elbow down is shown).

## Prismatic Joint no link offset

Joint 2 is prismatic.

- Define Frames 0 and 1 as before: $x_{1}$ is collinear with the origin of Frame 0.
- Define Frame 2 such that $x_{2}$ is collinear with $x_{1}$

$$
T_{2}^{1}=\left[\begin{array}{ccc}
1 & 0 & 0 \\
0 & 1 & 0 \\
0 & 0 & 1
\end{array}\right]\left[\begin{array}{ccc}
1 & 0 & a_{2} \\
0 & 1 & 0 \\
0 & 0 & 1
\end{array}\right]=\left[\begin{array}{ccc}
1 & 0 & a_{2} \\
0 & 1 & 0 \\
0 & 0 & 1
\end{array}\right] \quad T_{2}^{1}=\left[\begin{array}{ccc}
1 & 0 & a_{2} \\
0 & 1 & 0 \\
0 & 0 & 1
\end{array}\right]
$$

## Prismatic Joints

Joint 2 is prismatic.

- Define Frames 0 and 1 as before: $x_{1}$ is collinear with the origin of Frame 0.
- Define Frame 2 such that $x_{2}$ is collinear with $x_{1}$
$T_{2}^{0}=\left[\begin{array}{ccc}\cos \theta_{1} & -\sin \theta_{1} & a_{1} \cos \theta_{1} \\ \sin \theta_{1} & \cos \theta_{1} & a_{1} \sin \theta_{1} \\ 0 & 0 & 1\end{array}\right]\left[\begin{array}{ccc}1 & 0 & a_{i} \\ 0 & 1 & 0 \\ 0 & 0 & 1\end{array}\right]=\left[\begin{array}{ccc}\cos \theta_{1} & -\sin \theta_{1} & \left(a_{1}+a_{2}\right) \cos \theta_{1} \\ \sin \theta_{1} & \cos \theta_{1} & \left(a_{1}+a_{2}\right) \sin \theta_{1} \\ 0 & 0 & 1\end{array}\right]$


## Inverse Kinematic Solution

As before, let $\left(o_{2, x}, o_{2, y}\right)$ denote the coordinates of the origin of Frame 2.
Since $\theta_{2}=0$, the orientation of the endeffector frame is completely determined by $\theta_{1}$ :

$$
\theta_{1}=\tan ^{-1} \frac{o_{2, y}}{o_{2, x}}
$$



Once we have a solution for $\theta_{1}$, we can directly solve the forward kinematic equations for $a_{2}$ :

$$
\left(a_{1}+a_{2}\right) \cos \theta_{1}=o_{2, x} \Rightarrow a_{2}=\frac{o_{2, x}-a_{1} \cos \theta_{1}}{\cos \theta_{1}}
$$

## Three-link RPR Arm



Joint 2 is prismatic.

$$
\begin{aligned}
T_{3}^{0} & =\left[\begin{array}{ccc}
\cos \theta_{1} & -\sin \theta_{1} & a_{1} \cos \theta_{1} \\
\sin \theta_{1} & \cos \theta_{1} & a_{1} \sin \theta_{1} \\
0 & 0 & 1
\end{array}\right]\left[\begin{array}{ccc}
1 & 0 & a_{2} \\
0 & 1 & 0 \\
0 & 0 & 1
\end{array}\right]\left[\begin{array}{ccc}
\cos \theta_{3} & -\sin \theta_{3} & a_{3} \cos \theta_{3} \\
\sin \theta_{3} & \cos \theta_{3} & a_{3} \sin \theta_{3} \\
0 & 0 & 1
\end{array}\right] \\
& =\left[\begin{array}{ccc}
\cos \left(\theta_{1}+\theta_{3}\right) & -\sin \left(\theta_{1}+\theta_{3}\right) & \left(a_{1}+a_{2}\right) \cos \theta_{1}+a_{3} \cos \left(\theta_{1}+\theta_{3}\right) \\
\sin \left(\theta_{1}+\theta_{3}\right) & \cos \left(\theta_{1}+\theta_{3}\right) & \left(a_{1}+a_{2}\right) \sin \theta_{1}+a_{3} \sin \left(\theta_{1}+\theta_{3}\right) \\
0 & 0 & 1
\end{array}\right]
\end{aligned}
$$

## Inverse Kinematic Solution



As with the three-link RRR arm,

- Parameterize the end effector frame by $\left(X_{e}, Y_{e}, \phi\right)$
- Use $\phi$ and $a_{3}$ to solve for $\left(o_{2, x}, o_{2, y}\right)$
- Solve for $\theta_{1}$ and $a_{2}$ using the two-link RP solution given above.
- $\boldsymbol{\theta}_{3}=\boldsymbol{\phi}-\boldsymbol{\theta}_{1}$


## Other Kinds of Robots

So far, we've looked only at simple planar arms with revolute joints.
Life becomes more complicated if

- We have prismatic joints with "link offsets"
- Robots are not planar (e.g., anthropomorphic arms)
- Robots have more joints than end-effector degrees of freedom (e.g., if a planar arm has four joint). Such robots are said to be redundant.

We'll need to find other ways to solve the inverse kinematics for such robots.

Let's see an example...

## Prismatic Joint with link offset

Joint 2 is prismatic.


$$
T_{2}^{1}=\left[\begin{array}{lll}
1 & 0 & 0 \\
0 & 1 & 0 \\
0 & 0 & 1
\end{array}\right]\left[\begin{array}{ccc}
1 & 0 & a_{2} \\
0 & 1 & l_{2} \\
0 & 0 & 1
\end{array}\right]=\left[\begin{array}{ccc}
1 & 0 & a_{2} \\
0 & 1 & l_{2} \\
0 & 0 & 1
\end{array}\right]
$$

## Prismatic Joint with link offset



## Three-link RPR Example

The geometric IKS solution promises to be painful, and this is still a simple planar arm.

$$
\begin{aligned}
T_{3}^{0} & =\left[\begin{array}{ccc}
\cos \theta_{1} & -\sin \theta_{1} & \left(a_{1}+a_{2}\right) \cos \theta_{1}-l_{2} \sin \theta_{1} \\
\sin \theta_{1} & \cos \theta_{1} & \left(a_{1}+a_{2}\right) \sin \theta_{1}+l_{2} \cos \theta_{1} \\
0 & 0 & 1
\end{array}\right]\left[\begin{array}{ccc}
\cos \theta_{3} & -\sin \theta_{3} & a_{3} \cos \theta_{3} \\
\sin \theta_{3} & \cos \theta_{3} & a_{3} \sin \theta_{3} \\
0 & 0 & 1
\end{array}\right] \\
& =\left[\begin{array}{ccc}
C_{13} & -S_{13} & \left(a_{1}+a_{2}\right) C_{1}-l_{2} S_{1}+a_{3} C_{13} \\
S_{13} & C_{13} & \left(a_{1}+a_{2}\right) S_{1}+l_{2} C_{1}+a_{3} S_{13} \\
0 & 0 & 1
\end{array}\right]
\end{aligned}
$$

## More General Robot Arms

Imagine finding the geometric IKS solutions for more complicated arms in 3D work spaces...




## Numerical Methods for Inverse Kinematics

- Denote the forward kinematic map by $F(q)$.
- $F(q)$ is a vector-valued function, each component gives one coordinate of the end-effector pose. If the end effector has $m$ degrees of freedom and the robot has $n$ joints, then

$$
F\left(q_{1}, \ldots, q_{n}\right)=\left[\begin{array}{c}
f_{1}(q) \\
\vdots \\
f_{m}(q)
\end{array}\right]
$$

- For the two-link arm, we would have

$$
F\left(\theta_{1}, \theta_{2}\right)=\left[\begin{array}{l}
x\left(\theta_{1}, \theta_{2}\right) \\
y\left(\theta_{1}, \theta_{2}\right)
\end{array}\right]
$$

- Denote by $x^{d}$ the desired value of the end-effector pose.
- For the two-link arm we would have

$$
x^{d}=\left[\begin{array}{c}
o_{2, x}^{d} \\
o_{2, y}^{d}
\end{array}\right]
$$

- The inverse kinematic solution is the vector $q^{d}$ such that $F\left(q^{d}\right)=x^{d}$.


## Iterative Methods

Because we cannot solve $\boldsymbol{F}\left(\boldsymbol{q}^{d}\right)=x^{d}$ for $\boldsymbol{q}^{d}$ (nonlinear equation with no closed-form solution), we take an iterative approach.

- Generate a sequence of values $\boldsymbol{q}^{0}, \boldsymbol{q}^{1}, \boldsymbol{q}^{2} \ldots$ until we find some $\boldsymbol{q}^{N}$ such that $\boldsymbol{F}\left(\boldsymbol{q}^{N}\right)$ is sufficiently close to $x^{d}$, i.e.,

$$
\left\|F\left(\boldsymbol{q}^{N}\right)-x^{d}\right\|<\epsilon
$$

- In general, at each iteration, we compute $q^{i+1}$ by

$$
q^{i+1}=q^{i}+\delta q
$$

- The only trick is to determine a good value for $\boldsymbol{\delta} \boldsymbol{q}$ at each iteration.


## Gradient Descent

- In general, for gradient descent methods, the goal is to minimize the value of some function $L(q)$ by choosing the updates according to $\delta \boldsymbol{q}=-\nabla \mathrm{L}\left(\mathbf{q}^{\mathrm{i}}\right)$ where the gradient is w.r.t. $\boldsymbol{q}$.
- To solve the inverse kinematics problem, we wish to minimize the error between $\boldsymbol{F}(\boldsymbol{q})$ and $\boldsymbol{x}^{\boldsymbol{d}}$.
- Let's define $L(q)$ in terms of the squared error:

$$
L(q)=\frac{1}{2}\left(F(q)-x^{d}\right)^{T}\left(F(q)-x^{d}\right)
$$

- For this problem, we define the iteration as

$$
q^{i+1}=q^{i}-\alpha_{i} \nabla \mathrm{~L}\left(\mathbf{q}^{\mathrm{i}}\right)
$$

- We can use the scalar $\alpha_{i}$ to determine the magnitude of the step size.
- But what exactly is $\nabla \mathrm{L}\left(\mathbf{q}^{\mathrm{i}}\right)$


## Calculating the Error Gradient $(n, m=2)$

$$
\begin{aligned}
L(q) & =\frac{1}{2}\left[\left(f_{1}(q)-x_{1}^{d}\right)\right. \\
& \left.\left.\quad f_{2}(q)-x_{2}^{d}\right)\right]\left[\begin{array}{l}
{\left[\begin{array}{l}
f_{1}(q)-x_{1}^{d} \\
f_{2}(q)-x_{2}^{d}
\end{array}\right]} \\
\end{array}=\frac{1}{2}\left(f_{1}(q)-x_{1}^{d}\right)^{2}+\frac{1}{2}\left(f_{2}(q)-x_{2}^{d}\right)^{2}\right.
\end{aligned}
$$

$$
\nabla L(q)=\left[\begin{array}{c}
\frac{\partial L}{\partial q_{1}} \\
\frac{\partial L}{\partial q_{2}}
\end{array}\right]
$$

$$
\frac{\partial}{\partial q_{1}} L(q)=\frac{\partial}{\partial q_{1}} \frac{1}{2}\left\{\left(f_{1}(q)-x_{1}^{d}\right)^{2}+\left(f_{2}(q)-x_{2}^{d}\right)^{2}\right\} \quad \frac{\partial}{\partial q_{2}} L(q)=\frac{\partial}{\partial q_{2}} \frac{1}{2}\left\{\left(f_{1}(q)-x_{1}^{d}\right)^{2}+\left(f_{2}(q)-x_{2}^{d}\right)^{2}\right\}
$$

$$
=\frac{\partial f_{1}}{\partial q_{1}}\left(f_{1}(q)-x_{1}^{d}\right)+\frac{\partial f_{2}}{\partial q_{1}}\left(f_{2}(q)-x_{2}^{d}\right)
$$

$$
=\frac{\partial f_{1}}{\partial q_{2}}\left(f_{1}(q)-x_{1}^{d}\right)+\frac{\partial f_{2}}{\partial q_{2}}\left(f_{2}(q)-x_{2}^{d}\right)
$$

Stack these to obtain the gradient.

## Calculating the Gradient (cont)

$\nabla L(q)=\left[\begin{array}{l}\frac{\partial L}{\partial q_{1}} \\ \frac{\partial L}{\partial q_{2}}\end{array}\right]=\left[\begin{array}{l}\frac{\partial f_{1}}{\partial q_{1}}\left(f_{1}(q)-x_{1}^{d}\right)+\frac{\partial f_{2}}{\partial q_{1}}\left(f_{2}(q)-x_{2}^{d}\right) \\ \frac{\partial f_{1}}{\partial q_{2}}\left(f_{1}(q)-x_{1}^{d}\right)+\frac{\partial f_{2}}{\partial q_{2}}\left(f_{2}(q)-x_{2}^{d}\right)\end{array}\right]$
We can write this as a matrix equation

$$
\begin{aligned}
& =\left[\begin{array}{ll}
\frac{\partial f_{1}}{\partial q_{1}} & \frac{\partial f_{2}}{\partial q_{1}} \\
\frac{\partial f_{1}}{\partial q_{2}} & \frac{\partial f_{2}}{\partial q_{2}}
\end{array}\right]\left[\begin{array}{l}
f_{1}(q)-x_{1}^{d} \\
f_{2}(q)-x_{2}^{d}
\end{array}\right] \\
& =J^{T}(\boldsymbol{q})\left(\boldsymbol{f}(\boldsymbol{q})-\boldsymbol{x}^{\boldsymbol{d}}\right)
\end{aligned}
$$

$$
\text { And we recognize that }\left[\begin{array}{ll}
\frac{\partial f_{1}}{\partial q_{1}} & \frac{\partial f_{2}}{\partial q_{1}} \\
\frac{\partial f_{1}}{\partial q_{2}} & \frac{\partial f_{2}}{\partial q_{2}}
\end{array}\right]=J^{T}(q) \text {, the arm }
$$

Jacobian,

$$
\text { and that }\left[\begin{array}{l}
f_{1}(q)-x_{1}^{d} \\
f_{2}(q)-x_{2}^{d}
\end{array}\right]=\left(f(q)-x^{d}\right) \text { ! }
$$

And our gradient descent update becomes

$$
\boldsymbol{q}^{\boldsymbol{i + 1}}=q^{i}-\alpha_{i} \nabla \mathrm{~L}\left(\mathrm{q}^{\mathrm{i}}\right)=\boldsymbol{q}^{\boldsymbol{i}}+\boldsymbol{\alpha}_{\boldsymbol{i}} \boldsymbol{J}^{\boldsymbol{T}}\left(\boldsymbol{q}^{i}\right)\left(\boldsymbol{x}^{\boldsymbol{d}}-\boldsymbol{f}\left(\boldsymbol{q}^{\boldsymbol{i}}\right)\right)
$$

## Inverse Jacobian Method

Let's take a look at the Taylor series expansion for the forward kinematic map around the i-th iteratate:

$$
F\left(q^{i}+\delta q^{i}\right)=F\left(q^{i}\right)+J\left(q^{i}\right) \delta q^{i}+\text { h.o.t. }
$$

Here, h.o.t. refers to higher order terms (these go to zero quickly as $\delta q^{i} \rightarrow 0$ ). If we neglect the higher order terms, the ideal choice for $\delta q^{i}$ would be

$$
\delta q^{i}=q^{d}-q^{i}
$$

in which case we would step to the goal configuration in a single step!
In this case,

$$
F\left(q^{i}+\delta q^{i}\right)-F\left(q^{i}\right) \approx J\left(q^{i}\right) \delta q^{i}
$$

or

$$
x^{d}-F\left(q^{i}\right) \approx J\left(q^{i}\right) \delta q^{i}
$$

and this leads to the update

$$
\delta q^{i}=J^{-1}\left(q^{i}\right)\left(x^{d}-F\left(q^{i}\right)\right)
$$

Remember - we don't know $q^{d}$.
Happily, $q^{d}$ does not appear on the r.h.s!

## What about redundant arms...

- For the two-link arm, we can position the end-effector origin anywhere in the arm's workspace: two inputs ( $\theta_{1}, \theta_{2}$ ) and two "outputs" ( $X_{e}, Y_{e}$ ).
- For the three-link arm, we can position the end-effector origin anywhere in the arm's workspace, and we can choose the orientation of the frame: three inputs $\left(\theta_{1}, \theta_{2}, \theta_{3}\right)$ and three "outputs" ( $\left.X_{e}, Y_{e}, \phi\right)$.
- Suppose we had a four-link arm?
- Infinitely may ways to achieve a desired end-effector configuration $\left(X_{e}, Y_{e}, \phi\right)$.


## The case for $n>m$

In this case, there are "extra" joints:

$$
F\left(q^{i}+\delta q^{i}\right)-F\left(q^{i}\right) \approx J\left(q^{i}\right) \delta q^{i}
$$

If we write this out in detail, we see

$$
\left[\begin{array}{c}
f_{1}\left(q^{i}+\delta q^{i}\right) \\
\vdots \\
f_{m}\left(q^{i}+\delta q^{i}\right)
\end{array}\right]-\left[\begin{array}{c}
f_{1}\left(q^{i}\right) \\
\vdots \\
f_{m}\left(q^{i}\right)
\end{array}\right] \approx\left[\begin{array}{ccc}
\frac{\partial f_{1}}{\partial q_{1}} & \cdots & \frac{\partial f_{1}}{\partial q_{n}} \\
\frac{\partial f_{m}}{\partial q_{1}} & \cdots & \frac{\partial f_{m}}{\partial q_{n}}
\end{array}\right]\left[\begin{array}{c}
\delta q_{1}^{i} \\
\vdots \\
\delta q_{n}^{i}
\end{array}\right]
$$

Since J is not square, we can't invert it.

- Suppose $\operatorname{rank}(J)=m$.
- Then $J J^{T} \in \mathbb{R}^{m \times m}$ and $\operatorname{rank}\left(J J^{T}\right)=m$
- Then $\left(J J^{T}\right)^{-1}$ exists

What can we do with this...

## Pseudoinverses

- Define $J^{+}=J^{T}\left(J J^{T}\right)^{-1}$
- Suppose we let $\delta q^{i}=J^{+}\left\{F\left(q^{i}+\delta q^{i}\right)-F\left(q^{i}\right)\right\}$

Then

$$
J \delta q^{i}=J J^{+}\left\{F\left(q^{i}+\delta q^{i}\right)-F\left(q^{i}\right)\right\}=F\left(q^{i}+\delta q^{i}\right)-F\left(q^{i}\right)
$$

- In other words,

$$
J^{+}\left\{F\left(q^{i}+\delta q^{i}\right)-F\left(q^{i}\right)\right\}=\delta q^{i}
$$

is a solution to the equation

$$
F\left(q^{i}+\delta q^{i}\right)-F\left(q^{i}\right)=J\left(q^{i}\right) \delta q^{i}
$$

- We can use this to define our update law:

$$
\delta q^{i}=J^{+}\left\{F\left(q^{i}+\delta q^{i}\right)-F\left(q^{i}\right)\right\}
$$

$\cdot J^{+}$is called a pseudoinverse.

